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# Coherence of thermal transitions in poly(*N*-vinyl pyrrolidone)–poly(ethylene glycol) compatible blends 1. Interrelations among the temperatures of melting, maximum cold crystallization rate and glass transition

M.M. Feldstein\*, G.A. Shandryuk, S.A. Kuptsov, N.A. Plate´

*A.V. Topchiev Institute for Petrochemical Synthesis, Russian Academy of Sciences, 29 Leninsky prospekt, 117912, Moscow, Russian Federation*

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## **Abstract**

The phase behaviour of blends of high-molecular weight poly(*N*-vinyl pyrrolidone) (PVP) with short-chain poly(ethylene glycol) (PEG) of  $M_w = 400$ , prepared by drying their solutions in a common solvent (ethyl alcohol), was studied using DSC. Upon heating of cool-quenched samples a single glass transition was observed, followed by an exotherm corresponding to cold crystallization of excess PEG, a melting endotherm, and an endotherm corresponding to vaporization of absorbed water. The temperatures of glass transition (*T*g), PEG cold crystallization (*T<sub>c</sub>*), and melting (*T<sub>m</sub>*), along with the change in heat capacity ( $\Delta C_p$ ) between the polymer's glassy and rubbery states at  $T_g$ , vary with blend composition and hydration. As a result the  $T_g/T_m$ ,  $T_c/T_m$  and  $T_c/T_g$  ratios for PVP–PEG blends are functions of composition. PVP–PEG compatibility is due to H-bonding of PEG terminal hydroxyls to the carbonyls in the PVP repeating units. Large negative deviations of  $T_g$  values from the calculated weight averages, found mainly for PVP-overloaded blends, signify strong PVP–PEG interaction and free volume formation.  $© 2000$  Elsevier Science Ltd. All rights reserved.

*Keywords*: Poly(*N*-vinyl pyrrolidone); Poly(ethylene glycol); Compatibility

# **1. Introduction**

Differential scanning calorimetry (DSC) thermograms, which plot heat capacity,  $C_p$ , versus temperature, provide a highly informative tool to study the phase behaviour of polymers. A range of important thermodynamic characteristics may be readily determined from DSC traces. Such characteristics are glass transition temperature  $(T<sub>g</sub>)$ , change in heat capacity at  $T_g (\Delta C_p)$ , heat and temperature of melting  $(\Delta H_m, T_m)$ , and enthalpy and temperature of polymer crystallization ( $\Delta H_c$ ,  $T_c$ ). For hydrophilic polymers containing water, taken up from the environment via vapour state or residual from material processing, the temperature and heat of water thermodesorption along with a content of sorbed water can also be measured on DSC heating scans.

A number of empirical rules have been established regarding the transition temperatures for various homopolymers [1–4]. Boyer [2] and Beamen [3] have demonstrated a linear correlation between  $T<sub>m</sub>$  and  $T<sub>g</sub>$ , with slope

Corresponding author. Fax: +7-095-230-2224.

and intercept related to chemical structure. Boyer classified polymers into two large groups as symmetrical (sterically nonrestricted, low cohesive energy density (CED) macromolecules with flexible chains) and unsymmetrical (sterically restricted, substituted polymers with stiff backbone and bulky side groups). However, in an extensive study of 132 polymers, subsequent investigations found no sharp division between the  $T_{g}/T_{m}$  ratios observed for symmetrical and unsymmetrical polymers. Thus  $T_{\phi}/T_{\rm m}$  (in K) quantity has been reported to range from 0.5 to 0.8 with an average value to be [1]:

$$
\frac{T_{\rm g}}{T_{\rm m}} = 0.666\tag{1}
$$

Moreover, this rule was also shown to be true for both organic and inorganic compounds [5].

Similarly, the ratio  $T_c/T_m$  (where  $T_c$  designates the temperature of maximum crystallization rate upon polymer cooling from melt) shows an almost constant value for a wide variety of metals, organic and inorganic substances, varying between 0.75 and 0.90. According to the Mandelkern's rule [6], the rate of polymer crystallization

*E-mail address:* mfeld@ips.ac.ru (M.M. Feldstein).

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is maximum at a temperature of about 0.9 of the melting point. Closer examination for polymers has found [7]:

$$
\frac{T_{\rm c}}{T_{\rm m}} = 0.82 - 0.83\tag{2}
$$

In addition, the  $T_c$  relates to  $T_g$  by the equation [1,8,9]:

$$
\frac{T_c}{T_g} = 1.20 - 1.33\tag{3}
$$

Following Boyer [2], the minimum temperature required for crystallization,  $T_c$  (°C), is also a function of  $T_g$ . For a wide range of polymers, the points fall on a straight line, defined by the equation:

$$
T_c(^{\circ}\text{C}) = 32.5 + 1.125 T_g(^{\circ}\text{C})
$$
 (4)

More general empirical relations between  $T_{\rm g}$ ,  $T_{\rm m}$  and  $T_{\rm c}$ include consideration of all three transition temperatures [5]:

$$
T_{\rm c} = \frac{(T_{\rm g} + T_{\rm m})}{2} \text{ or } T_{\rm c} - T_{\rm g} = T_{\rm m} - T_{\rm c}
$$
 (5)

Although the rules  $(1)$ – $(5)$  are totally empirical, they invoke fundamental structure–property relationships, which have begun only recently to find quantitative comprehension. The  $T_g$  is defined as a function of coordination number, *z*, and total interaction energy,  $\langle Do \rangle$ , of the atoms forming the polymer segment [10]:

$$
T_{\rm g} = 0.445 \frac{z \langle Do \rangle}{\rm R} \tag{6}
$$

where R is the universal gas constant. Since  $z$  is in inverse proportion to free volume (Eq.  $(6)$ ) outlines the energy– volume ratio, related to CED. This definition of  $T<sub>g</sub>$  is in agreement with the results of recent work [11], dealing with the calculations of  $T_g$  values for the series of aliphatic acrylate and methacrylate polymers on the basis of energy– volume–mass (EVM) model.

Like  $T_{\rm g}$ , the melting points of polymers are greater the greater the CED, and the greater the chain stiffness, i.e. the smaller the free volume [2]. The physical meaning of the constancy of the  $T_{g}/T_{m}$  ratio (1) is that the molecular packing coefficients of crystalline polymers at  $T<sub>m</sub>$  approximate those of amorphous polymers at  $T_g$ , that is the fusion of polymer crystals and transition of amorphous polymer glasses into viscoelastic state occurs upon achieving the same values of the fraction of full free volume [12], equal to  $1 - (T_g/T_m) = 1 - 0.666 = 0.333$ .

Based on an iso-volume state model and a crystallization theory, the relationships between  $T_{\text{m}}$ ,  $T_{\text{c}}$  and  $T_{\text{g}}$  are formulated by Okui [1] as follows:

$$
\frac{T_g}{T_m} = \frac{C - (\alpha_c/\alpha_g)}{C + 1},\tag{7}
$$

$$
\frac{T_c}{T_m} = \frac{C}{C+1},\tag{8}
$$

$$
\frac{T_{\rm m} - T_{\rm c}}{T_{\rm c} - T_{\rm g}} = \frac{\alpha_{\rm g}}{\alpha_{\rm c}},\tag{9}
$$

$$
C = \sqrt{1 + (\Delta E/K)}\tag{10}
$$

where  $\alpha_{\rm g}$ ,  $\alpha_{\rm c}$  are the thermal expansion coefficients of glassy and crystalline polymers, respectively;  $\Delta E$  is an activation energy for polymer segment migration through the nucleus–melt interface and *K* is the nucleation parameter associated with the mean surface energy,  $\delta$ , and the heat of fusion,  $\Delta H_m$ :

$$
K = \frac{n\delta^2}{\Delta H_{\rm m}}, \qquad \delta = \sqrt{h_0 \sigma_{\rm e} \sigma_{\rm m}}
$$
(11)

In Eq. (11), *n* is a nucleation mode parameter,  $h_0$  the thickness of the depositing growth layer, and  $\sigma_{\rm e}$  and  $\sigma_{\rm m}$  the end and the lateral surface energies, respectively. In general, *C* varies from 3 to 9 with a mean value of about 5. The average  $\Delta E/K$  magnitude is around 23 for most polymers [1].

 $\Delta E$  may be compared with the activation energy for viscous flow or self-diffusion of macromolecular segments, which are the characteristics of chain mobility.  $\Delta H_{\text{m}}$  is a measure of intermolecular forces and  $\delta$  affects the degree of polymer crystallinity [1]. All three parameters may be expressed in terms of cohesive energy, chain stiffness and geometry, which contribute to the free volume. Both  $T<sub>g</sub>$  and *T*<sup>m</sup> are shown to invoke the same fundamental properties of polymers [2,10–12].

The numerical quantities outlined by Eqs.  $(1)$ – $(5)$  and (7)–(9), hold only for individual homopolymers and, in general, are inapplicable for copolymers or polymer blends [10,12]. The view taken in this series of papers is that a blend of compatible polymers can behave like an individual homopolymer, demonstrating the compositional dependence of the quantities defined by Eqs.  $(1)$ – $(10)$ . In this work, we follow terminology defined by Paul and Newman [13]: a blend is considered compatible if mixed polymers are capable to exhibit a total miscibility on a molecular scale. Miscibility on the molecular scale is not necessarily random within the entire range of composition and temperatures: interactions between similar or different macromolecules may lead to a certain amount of clustering or other nonrandom arrangement of polymer segments. Thus, the crystallization of polymer in the blend is not treated as a sign of incompatibility if the crystallizable polymer remains at least partially amorphous and miscible with the other component in the amorphous state below  $T_{\text{m}}$  [13].

In many blend systems, a homogeneous phase is obtained because of the existence of specific favourable interactions between different polymer components, which allow mixing on a molecular scale. One such favourable interaction is H-bonding that has been reported for many polymer blends [14–17]. Polymers containing ternary amide groups, such as poly(*N*-vinyl pyrrolidone) (PVP), are potentially good proton acceptors due to the basic nature of the functional groups [18]. In a similar way, another good proton-accepting polymer is poly(ethylene oxide) (PEO) [19], and therefore it comes as no surprise that PVP and PEO are reported to form only incompatible binary blends [20], while in ternary blends with such proton-donating compatibilizers as poly- (acrylic acid) and poly(methacrylic acid) they display compatibility [17]. As the more potent competitor in H-bonding with polyacids, in doing so the PVP supersedes PEO in forming hydrogen-complex with the protondonating polymer [21].

In contrast to high molecular weight PEO, short-chain poly(ethylene glycol) represents a symmetric telechelic polymer with controlled molecular weight and narrow molecular weight distribution, carrying two proton-donating hydroxyl groups at the chain ends. The term "PEG" is used in this work to define the oligomeric starting polymers containing terminal hydroxyls, while the "PEO" is employed to emphasize that the contribution of end-chain hydroxyls to miscibility is considered to be negligible due to higher molecular weight and lower end-group concentration. PVP compatibility with short-chain PEG follows directly from the well-known observation of unlimited PVP solubility in liquid PEG [22]. In this way, the PVP–PEG compatible blends may be treated as the solutions of high molecular weight PVP in liquid oligo(ethylene glycol).

PVP–PEG compatibility in blends has been determined by FTIR spectroscopy to be due to H-bonded interactions between the hydrogen atom of PEG terminal groups and electronegative oxygen atom in the carbonyl groups of the monomer units of the comparatively longer PVP chains [23,24]. High molecular weight PVP forms compatible blends only with short-chain PEGs [25], ranging in molecular weight from 200 to 600 g mol<sup>-1</sup>. PEGs of higher molecular weight are miscible with amorphous PVP in the melt, but these do not form homogeneous amorphous mixture with PVP below  $T_m$ . This is indicative of PVP incompatibility with long-chain PEG and PEO on the molecular scale, that is shown by direct optical microinterference measurement of PVP–PEG spontaneous mixing [26], and confirmed with DSC and wide angle X-ray scattering (WAXS) [27]. Microphase separation in PVP blends with PEO  $(M_w = 32\,000)$  has been also demonstrated by Cesteros et al. [20] with DSC and DMTA. These findings are explicable, because with the rise in PEG molecular weight the end-group contribution to compatibility becomes negligible.

H-bonding of PEG terminal hydroxyls to carbonyls in the PVP repeat units results in a high degree of the orientation of the PEG macromolecules with respect to longer PVP chains, as has been established with WAXS and rheological techniques [23,28]. The PVP–PEG complex is akin to those formed by H-bonds between the units of polar polymers and complementary end-groups of a range of oligomers, such as phenolic surfactants [29,30], but the former is distinguishable by its capability to form simultaneously two strong H-bonds through both PEG terminal groups. The

mechanism of PVP–PEG H-bonding implies the crosslinking of PVP macromolecules by relatively long and flexible PEG chains. The resulting network structure of PVP–PEG complex reveals a range of unique physical properties, including a high elasticity, similar to that of lightly cross-linked rubbers [23,31], and pressure-sensitive character of adhesion toward various substrates [32–34]. With conventional compounding methods, polymer blends, which possess high tack, do so because their parent polymers are also tacky. This is not the case for the present PVP–PEG adhesive, where adhesion has been found to arise within a narrow range of PEG concentration [32]. Tailoring adhesive properties by blending two non-adhesive polymers, PVP and PEG, provides an indirect evidence of their compatibility and the formation of new, distinguishable and independent supramolecular species. Indeed, the PVP–PEG blends have been defined as a complex of stoichiometric composition, in which approximately 30% of PVP repeat units are bonded to PEG terminal hydroxyls and nearly 15 PEG chains are associated on average with 100 PVP units [28].

This present series of papers report on the phase behaviour of the PVP–PEG blends as examined by DSC of cool-quenched samples under heating. In contrast to  $T_c$ definition given above, the  $T_c$  is here defined as the temperature at which the rate of crystallization is maximum and coincides with a peak temperature of cold crystallization under heating through  $T<sub>g</sub>$  of a blend, produced by quenchcooling from the melt. The first paper considers the correlation between temperatures of glass transition, PEG cold crystallization and melting. The second paper presents the thermodynamic analysis of  $T_c$  dependence on  $T_g$ . The third paper describes the impact of sorbed water upon relaxation and phase transitions.

#### **2. Experimental**

PVP (Kollidon K-90),  $M_w = 1000000 \text{ g mol}^{-1}$ , and PEG (Lutrol E-400),  $M_w = 400 \text{ g mol}^{-1}$ , were obtained from BASF and used as-received. Both polymers are hygroscopic and the degree of their hydration, evaluated by the weight loss under drying at  $105^{\circ}$ C, is taken into account to prepare physical blends spanning the entire range of compositions. Depending on the relative humidity (RH) of the surrounding atmosphere, PVP hydration ranged from 6 to 8 wt%, while the PEG contained less sorbed water  $(0-1 \text{ wt\%})$ .

The blends, referred throughout this work as "fresh", were produced by dissolving PVP and PEG in common solvent (ethyl alcohol) followed by the solvent evaporation at ambient temperature and RH until a constant weight was attained. Solvent removal was verified by FTIR spectroscopy, observing the lack of methylene group stretching vibrations at 2974 and 1378 cm<sup> $-1$ </sup> in the IR spectrum. The hydration of freshly prepared PVP–PEG blends averaged





 $7.2 \pm 2.2$  wt%. From this point on the term "hydrogels" is used to designate equilibrium hydrated PVP–PEG blends.

Since our earlier obtained data on PVP–PEG H-bonding [24] and water vapour sorption [35] provide grounds for considering sorbed water as the third component of the PVP–PEG complex, our concern here is confined mainly at the DSC study of PVP–PEG hydrogel phase behaviour. However, being a common solvent for both PVP and PEG, water is supposed to serve as a compatibilizer in PVP–PEG blends. To test this hypothesis, so-called "dry" blends were prepared by drying hydrogels at  $105^{\circ}$ C until termination of weight loss. Strictly speaking, efforts to prepare perfectly anhydrous samples were unsuccessful because, due to the high affinity for water vapour, dry samples contained  $0.7 \pm 0.2$  wt% of moisture sorbed in the course of sample handling and weighing. Dry blends were stored until use over  $P_2O_5$ , while hydrogels were either exposed to atmospheric humidity or equilibrated in a desiccator over aqueous  $H_2SO_4$  solution of controlled density  $(1.335 \text{ g cm}^{-3})$  which maintained the required  $RH = 50\%$ .

The samples of PVP–PEG blends were subjected to several different DSC experiments to assess the effects of various test procedures (heating rates) and storage conditions (time and RH) on the transition temperatures and enthalpy changes occurring in each system. Approximately 5–15 mg samples of each composition were sealed in standard aluminium pans with lids pierced with a sharp pin to enable the evaporation of sorbed water. The samples were analysed under a dry argon purge  $(50 \text{ ml min}^{-1})$  in a Mettler TA 4000/DSC 30 DSC, calibrated for temperature and heat flow using indium and gallium ultrapure standards. All reported values are the average of replicate experiments varying less than 1–2%. All the samples were reweighed after scanning and weight loss was registered.

In the DSC apparatus, samples were first cooled with liquid nitrogen from ambient temperature to  $-100^{\circ}$ C over 2–3 min and then heated at a rate of  $20^{\circ}$ C min<sup>-1</sup> (unless otherwise specified) to  $200^{\circ}$ C. Upon heating, a heat capacity jump followed by single exotherm coupled with symmetric endotherm, and high-temperature endotherms were normally observed for PVP–PEG blends. These four transitions were, respectively, attributed to the glass transition, PEG cold crystallization, melting, and water thermodesorption, as explained in the following section of this paper.  $T_{\varphi}$ s were recorded at the half-height of the corresponding heat capacity jumps, whereas the  $T_c$ ,  $T_m$  and water thermodesorption temperatures were taken as the relevant peak temperatures. Heats of isolated peaks were computed by constructing linear baselines from the peak onset to completion and numerically integrated with appropriate software supplied by Mettler. For overlapping peaks, the baselines were constructed from the onset of the first peak to the completion of the second peak and the integration was carried out separately, below and above the valley between

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Polymer	Sorbed water (wt%)		Weight loss (% after scanning)
	<b>DSC</b>	<b>DTG</b>	
PVP <sup>a</sup>	6.0	4.18	5.2
$PVP-PEG (36 wt%)$	7.3	7.78	8.1
$PVP-PEG (69 wt%)$	1.9	1.86	2.0
PEG	0.2	0.44	0.4

The hydration of PVP, PEG and their blends tested with DSC and DTG compared to the weight loss of specimens after scanning

<sup>a</sup> As obtained from manufacturer. The hydration of PVP film dried from solution in ethyl alcohol vary from 10 to 12%.

the peaks. Due to dissimilarity in sample weights, the magnitudes of heat flow, presented in DSC traces throughout this work have been reduced everywhere to a reference weight of 10 mg.

# **3. Results and discussion**

Table 2

Characteristics of the phase transitions reported in this work relate to a scanning rate of  $20^{\circ}$ C min<sup>-1</sup>. The majority of the transitions are heating-rate dependent and the rate effects upon  $T_g$ ,  $T_c$ ,  $T_m$ ,  $\Delta H_c$ ,  $\Delta H_m$ , and the temperature of water desorption,  $T<sub>d</sub>$ , are displayed in Table 1. To generate the table, blends were heated at the rates of 10, 20, 30 and  $40^{\circ}$ C min<sup>-1</sup>, and isothermal characteristics were estimated by extrapolating relevant linear relationships to zero rate. Linear regression parameters, *R* and *p*, show statistically insignificant rate-dependence for  $T_{\text{m}}$ , while the  $T_{\text{g}}$ ,  $T_{\text{c}}$ ,  $T_{\text{d}}$ ,  $\Delta H_c$  and  $\Delta H_m$ -relationships are fairly fitted with linear graphs being plotted against both the rate of heating and the square root of this rate. The apparent  $T<sub>m</sub>$  independence on heating rate is most likely due to dominant depression of blend crystallinity by hydration.  $T<sub>m</sub>$  and the heat of water thermodesorption  $(\Delta H_d)$  are much more affected by blend hydration than PEG content, as disclosed in the third article



Fig. 1. DSC heating thermograms for PVP–PEG blends over the range of PEG/PVP compositions (in wt%): (1) 100% PVP; (2) 36% PEG; (3) 52.9% PEG; (4) 69.2% PEG; (5) 84.9% PEG; (6) 100% PEG. Storage time and conditions: three months at  $RH = 50\%$ , 25°C.

of this series. Reasons for the heating rate dependence of the area under the melting endotherm,  $\Delta H_{\text{m}}$ , are considered in the second paper of this series. Measured at the heating rate of 20 $\degree$ C min<sup>-1</sup> the  $T_g$  and  $T_c$  values are on average 4 and  $6\degree$ C higher than their respective magnitudes found by extrapolation to zero rate. With the increasing PEG concentration the difference between measured and extrapolated  $T_g$ and  $T_m$  values tends to decrease. The  $\Delta H_m$  and  $\Delta H_c$  vary linearly with the content of crystalline PEG and deviations from this rule, displayed in Table 1, are caused by a hydration effect, as shown in the third paper.

The loss in weight of a specimen after thermal scanning is not always indicative of its hydration degree but may also be due to partial material decomposition at high temperatures. The DSC technique is reported to provide an alternative tool to determine the amount of water sorbed by hydrophilic polymers [36]. To accomplish this, the heat of water thermodesorption must match that of bulk water evaporation. Polymer hydration may then be measured through the enthalpy change associated with water desorption, dividing the obtained dehydration enthalpy by the reference value of water vaporization,  $\Delta H = 2255 \text{ J g}^{-1}$ . This measurement is justified if specific water sorption in PVP–PEG hydrogels and polymer components does not contribute appreciably in the enthalpy of water thermodesorpion, i.e. the energy of water H-bonding to polymer units is comparable with the energy of H-bonds, formed between water molecules in clusters. To be certain of the validity of using the reference  $\Delta H$  value, the amounts of water sorbed with PVP–PEG blends and initial polymers, as obtained from DSC heating thermograms, were compared to those measured by differential thermogravimetry (DTG) technique. The identical blends were simultaneously heated with a rate of  $10^{\circ}$ C min<sup>-1</sup> from 25 to 200 $^{\circ}$ C in the DSC apparatus and in a Mettler TG 50 DTG analyser. The results are presented in Table 2 along with the weight loss of specimens after scans. Both measurements give similar results, confirming the validity of using DSC for evaluation of water content in PVP and PVP–PEG blends. The relative inaccuracy of DSC technique in assessing water sorption with pure PEG is the result, most likely, of very low water sorption.

DSC traces for PVP, PEG and their blends are shown in Fig. 1. The scan of unblended PVP reveals a broad symmetric endotherm of water thermodesorption at  $116^{\circ}$ C



Fig. 2. The  $T<sub>g</sub>$  of dry PVP–PEG blends as a function of PEG weight fraction. Points—experimental values, lines—the relationships described by the Fox and the Gordon–Taylor  $(k = 0.55)$  equations.

followed by a heat capacity jump at the glass transition  $(T_g = 178^{\circ}\text{C}, \ \ \Delta C_p = 0.27 \text{ J g}^{-1} \text{ K}^{-1}$ . The amount of water sorbed with PVP, determined as a ratio of water thermodesorption enthalpy to the reference value of bulk water vaporization heat, is found to be 7.3 wt%. This value is in close agreement with the loss in the sample weight after scanning (6.6 wt%). Immediate rescanning of the specimen using the same heating program shows neither endotherm nor weight loss, while the parameters of PVP glass transition remain practically intact. No freezing water is detected under PVP hydration even if 20% water sorption is attained.

The DSC heating curve of unblended PEG following its exposure to water vapour at  $RH = 50\%$  (Fig. 1) displays glass transition at  $-70^{\circ}$ C,  $\Delta C_p = 0.32$  J g<sup>-1</sup> K<sup>-1</sup> and PEG fusion endotherm at  $T_m = 6^\circ\text{C}$ ,  $\Delta H_m = 118.4 \text{ J g}^{-1}$ . The PEG is nearly six times less hydrated than PVP, and the broad peak of water desorption is shifted to  $127^{\circ}$ C. As the content of sorbed water increases to 11.3%,  $\Delta C_p$ achieves the value 1.17 J  $g^{-1} K^{-1}$  and the glass transition in PEG becomes more conspicuous  $(T_{\rm g} = -71^{\circ}\text{C})$ . Again, no freezing water is found as the PEG hydration is so high as 17 wt%, corresponding to maximum bound water content reported to be 2.7 water molecules per PEG unit [37,38].

Since the PVP and PEG glass transition temperatures differ by a value of about 250°C, the peculiarities of  $T<sub>g</sub>$ compositional behaviour in their blends are easily discernible (Fig. 1). A single composition dependent  $T_g$ , intermediate between those of the pure components, is an unambiguous criterion of the PVP–PEG compatibility. In the analysis of PVP–PEG miscibility, not only the number of glass transitions and their location on the temperature scale, but also their breadth were considered. Both the onset and endpoint of the glass transition in PVP–PEG compatible blends were recorded: the onset  $T_g$  was evaluated at the intersection of the pre-event baseline and a line drawn tangentially to the inflection point, while the endpoint was evaluated at the intersection of the inflection point tangent and the baseline established after the thermal event. A large transition width in PVP-overloaded blends, defined as the difference between the endpoint and onset  $T_{g}$ s and as a difference between inflection point and midpoint of glass transition, implies a broad spectrum of segment mobilities due to local heterogeneties caused by locally different degrees of PVP–PEG interaction. As PEG concentration in blend increases, both the cooperativeness of glass transition and  $\Delta C$ <sub>p</sub> become greater (Fig. 1).

Upon heating quench-cooled samples through  $T_g$ , the PVP-overloaded blends exhibit only the uppermost endotherms of water thermodesorption, while for PEGoverloaded mixtures the exotherms of PEG cold crystallization coupled with the PEG melting endotherms appear within intervening temperature ranges (Fig. 1). The water thermodesorption endotherms occur only in the scans of first heating and disappear at immediate rescanning. The enthalpy of water thermodesorption correlates with weight loss of the specimen after heating. (The effect of composition on this endotherm is considered in the third paper of this series.) By contrast, the enthalpies of the intervening exotherm and endotherm relate directly to PEG concentration in the blends and are only slightly affected by hydration. Although free water fusion in hydrogels occurs in the same temperature region, in PVP–PEG blends the thermal peaks characterize the PEG phase behaviour. Indeed, the peaks are equally inherent in both dry and hydrated blends, occurring both in the first and in the subsequent scans. These features argue unequivocally in favour of assigning the relevant peaks to excess PEG cold crystallization and fusion.

Since only the glass transition is featured without exception for all the dry and hydrated, amorphous and crystalline PVP–PEG blends, we start our analysis with the consideration of  $T_g$  compositional behaviour. While the scrutinity of cold crystallization–melting and vaporization peaks can be augmented to provide intrinsic information about the states of PEG and water in miscible blends, only the glass transition can be used to signify the mechanism of PVP–PEG interaction over entire composition range.

A variety of equations have been proposed to express the *T*g-composition dependence in miscible polymer blends and plasticized systems, as reviewed by Aubin and Prud'homme [39]. In general, it is observed that  $T_g$  varies monotonically as a function of composition and the difference between measured  $T_g$  values and those predicted with relevant equations is usually considered as a measure of the strength of interactions between molecules of the involved components. When blends are formed from strongly interacting pairs, specific bonds bridge long sequences of repeat units of complementary chains, decreasing free volume between them and increasing packing density and the energy of cohesion. In full agreement with Eq. (6), the  $T<sub>g</sub>$  of the blends in this instance are generally much higher than the values calculated as weight averages of the component polymers. Large positive deviations of  $T_g$  values from those predicted with the Fox [40] and Gordon–Taylor [41]



Fig. 3. The  $T_g$  of dry and freshly prepared (hydrated) PVP–PEG blends plotted against the number of PEG chains per one PVP unit available in the blend.

equations now appear to be a common occurrence in interpolymeric H-bonded and cation–anion complexes [42–45]. Large positive deviations are typical of PVP blends with hydroxyl-containing epoxy resin [46]. As regards to PVP H-bonded complexes with other hydroxyl containing polymers, such as polyvinyl alchohol [47,48], poly(hydroxyethyl methacrylate), poly(hydroxypropyl methacrylate) [49], and poly(*N*-phenyl-2-hydroxytrimethylene amine) [50], as well as for PVP miscible blends with polysulfone [51], poly(amide enaminonitile) [17], and phenoxy resin [52], the plots of  $T_g$  against the blend composition have been shown to follow fairly well the Fox, Gordon–Taylor, and Kwei [53] equations [47].

In contrast, H-bonding in PVP–PEG compatible blends results in large negative deviations from the simple rules of mixing expressed by the Fox and Gordon–Taylor equations (Fig. 2). With only 36 wt% of liquid PEG-400 added to glassy PVP, the blend  $T_g$  drops dramatically over 220°C, approaching the  $T_g$  value found for pure PEG. Following PVP mixing with plasticizer demonstrates a gradual levelling



Fig. 4. The fractional free volume,  $f$ , calculated with Eq. (12) from  $T<sub>g</sub>$ values, as a function of PVP–PEG blends composition at 20°C.

off the response in blend  $T<sub>g</sub>$  (Figs. 2 and 3). Within the frameworks of the classical polymer plasticization–dissolution concept the first stage of the PVP–PEG mixing may be identified as PVP plasticization, whereas the second stage consists in gradual dissolution of the plasticized polymer in excess plasticizer [54]. The driving force for PEGs plasticizing action is the H-bonding of hydroxyl groups at the ends of PEG short chains to carbonyls in PVP repeat units as has been established with FTIR spectroscopy [23]. The first stage represents, therefore, PVP–PEG complexation and the boundary between the two stages (i.e. the blend containing 36 wt% of PEG-400) corresponds to a PVP– PEG stoichiometric complex [25]. Upon achieving this PEG concentration in blend, the exothermic heat of H-bonds formation ceases to dominate the unfavourable change in PVP–PEG noncombinatorial interaction entropy, which has been shown to be negative [25]. By H-bonding PVP units through terminal hydroxyl groups, the PEG short and flexible chains behave as spacers increasing the free volume between PVP neighbour segments. The  $T<sub>g</sub>$  data in Figs. 2 and 3 enable in estimating the fraction of free volume, *f*, in PVP–PEG blends (Fig. 4) by the combination of the Doolittle and Williams–Landell–Ferry (WLF) equations [55]:

$$
\frac{1}{2.303} \left( \frac{1}{f} - \frac{1}{f_{g}} \right) = -\frac{Nf_{c}}{2.303f_{g}} \frac{T - T_{g}}{(f_{g}/\Delta\alpha) + T - T_{g}}
$$

$$
= -\frac{17.37(T - T_{g})}{51.6 + T - T_{g}}
$$
(12)

where  $f_g$  is the fractional free volume of a polymer at  $T_g$ ,  $\Delta \alpha$ is the change of thermal expansion coefficient of polymer at  $T_g$  ( $\Delta \alpha \approx 4.8 \times 10^{-4} \text{ K}^{-1}$ ),  $f_c$  is the critical fractional free volume required so that a segment may jump or move and *N* the number of moving units per segment. As was shown by Bueche,  $Nf_c \approx 1$ , Fox and Flory found  $f_g \approx 0.025$  for the majority of polymers [55]. Substitution of all the constants to Eq. (12) gives an expression ready to use.

The increase of free volume upon PVP–PEG mixing also appears from the compositional plot of another fundamental characteristic of glass transition,  $\Delta C_p$  (Fig. 5). As Tanaka has shown [56]

$$
\Delta C_{\rm p} = \left(\frac{\mathrm{d}}{\mathrm{d}T} \left[ \left\{ RT^2 \frac{\mathrm{d} \ln Z(T)}{\mathrm{d}T} \right\} / n \right] \right)_{\rm p}
$$

$$
- \left(\frac{\mathrm{d}}{\mathrm{d}T} \left[ RT^2 \frac{\mathrm{d} \ln V}{\mathrm{d}T} \right] \right)_{\rm p} + \left(\frac{\mathrm{d}H_0}{\mathrm{d}T} \right)_{\rm p} + 1.5R \qquad (13)
$$

where *Z*(*T*) is the steady-state conformational sum of a selected polymer chain at temperature *T*; *V* is the molar free volume per polymer segment; *n* the polymerization degree and  $H_0$  the molecular cohesive energy per one mole of polymer segments. The drop in  $\Delta C$ <sub>p</sub> at small PEG concentrations (Fig. 5) is the result of the dominant free volume contribution that has been found to arise at



Fig. 5. The change in heat capacity,  $\Delta C_p$ , between glassy and rubbery states plotted against the composition of freshly prepared and dry PVP–PEG blends.

PVP–PEG complex formation (Fig. 4). The subsequent growth in  $\Delta C_p$  is most likely a consequence of enhancing the energy of intermolecular interaction and H-bonding between macromolecules as PEG content in the blend increases. At minimum  $\Delta C$ <sub>p</sub> the energy of intermolecular cohesion within the PVP–PEG blend is specifically balanced by the fluctuation free volume. The observed reduction in  $\Delta C$ <sup>p</sup> for PEG-overloaded blends (Fig. 5) is due to their high crystallinity degree. It is known that  $\Delta C_p$ tends to decrease as crystallinity increases.

Although the  $T_g$  values of dry PVP–PEG blends are found to be somewhat higher than those for hydrated blends (Fig. 3), the sorbed water does not reveal itself as a compatibilizer in PVP–PEG system, because both dry and hydrated blends are compatible, demonstrating only one compositiondependent glass transition. This conclusion is presented in more quantitative terms in the third paper of this series.

Coherent behaviour of  $T_g$ ,  $T_m$  and  $T_c$  in PVP–PEG



Fig. 6. Effects of PVP–PEG blend composition and storage conditions on the relationship between glass transition and PEG melting temperatures: (1) aged blends; (2)  $RH = 50\%$ ; (3) freshly prepared blends.



Fig. 7. Effects of PVP–PEG blend composition and storage conditions on the relationship between PEG cold crystallization and melting temperatures upon heating quench-cooled samples through glass transition: (1) fresh blends; (2)  $RH = 50\%$ ; (3) February; (4) May.

compatible blends is illustrated in Figs. 6 and 7. In agreement with the rules outlined by Eqs. (1) and (2),  $T<sub>m</sub>$  is a linear function of  $T_g$  and  $T_g$ , however depending on the time and conditions of storage, the samples of different prehistory are characterized with noticeably different slopes and intercepts  $(R = 0.88 - 0.99)$ . Thus, freshly prepared blends exhibit lower  $T<sub>m</sub>$  values, but after being set aside for a period of up to one year, the blends show signs of long-term evolution or "ageing", which among other changes result in the increase in melting temperature. The melting temperature is affected by the perfection of the crystal structure and, consequently, structural relaxation processes underlie the accompanying changes with time in the phase transition temperatures of PVP–PEG blends.

In freshly prepared blends, produced by evaporation of common solvent at ambient temperatures, the PEG chains, attached to PVP units through H-bonding of terminal hydroxyls, are most likely randomly disordered in their positions along PVP chains. The self-assembly of several PEG chains into energetically more favourable ordered structures is supposed to be a driving force for the observed ageing. Similar self-assembly of ionic surfactant amphiphilic molecules into clusters is reported to accompany surfactant binding with oppositely charged polyelectrolytes [57]. The PEG chains orientation within such ordered structures has been recently shown by WAXS study of PVP–PEG blends to match closely that of PEG crystals [23]. It is obvious that the more the PEG chains being assembled into crystal-like structures, the higher the  $T<sub>m</sub>$ . On the other hand, sorbed water is capable of disrupting crystal perfection and depressing  $T<sub>m</sub>$ . Aged blends are appreciably drier than freshly prepared ones and exhibit higher *T*ms. The changes in blend hydration account for seasonal effects observed in phase behaviour of aged PVP–PEG blends. The blends tested in February after nine months of storage at comparatively lower RH are referred throughout



Fig. 8. The  $T_g / T_m$ ,  $T_c / T_g$  and  $T_c / T_m$  quantities as the functions of PEG concentration in freshly prepared blends.

this work to as "February", whereas those tested in May following the storage for a period of one year at higher RH of surrounding atmosphere are designated as "May" or simply "aged" blends. The analysis of seasonal effects is presented in the third paper of this series, in which the impact of hydration upon phase behaviour has been shown to dominate over the influence of storage time. For the purposes of the present discussion, the data sets obtained for different storage conditions are considered separately in relation to PEG content in blends, thereby minimising the effects of relative humidity.

Fig. 8 illustrates the dependence of  $T_g / T_m$ ,  $T_c / T_g$  and  $T_c / T_g$ *T*<sup>m</sup> quantities on the composition of freshly prepared PVP– PEG blends. These are explicit functions of blend composition, expressed in terms of PEG weight fraction. Being reduced to the most basic molecular level, the temperatures of phase transitions are interrelated through the structure of the polymer. In compatible polymer blends the phase behaviour results from polymer mixing at a molecular level and varies in accord with composition. As PEG concentration in the blend increases, the  $T_g/T_m$  value varies linearly from



Fig. 9. Effect of storage conditions on the compositional behaviour of  $T_g / T_m$ quantity in PVP–PEG compatible blends, freshly prepared and aged for three months at  $RH = 50\%$  and 25°C.

0.785 to 0.760, whereas other quantities decrease smoothly from  $T_c/T_m = 0.958$  to 0.849 and from  $T_c/T_g = 1.220$  to 1.111. Compared with reference average magnitudes found for a variety of individual polymers (Eqs.  $(1)$ – $(3)$ ) [1], in PVP–PEG blends the  $T_g/T_m$  and  $T_c/T_m$  values are at the upper edges of the relevant ranges, while the  $T_c / T_g$ value is at the lower boundary. It should be emphasized that the obtained values relating to the PEG cold crystallization peak temperatures upon heating quench cooled samples from glassy state,  $T_c$ , match closely the reference magnitudes obtained after cooling polymers from the melt. This finding implies that in PVP–PEG compatible blends the maximum cold crystallization rate temperature is in close agreement with the temperature of maximum crystallization rate. Since the quantity  $(1 - T_g/T_m)$  defines the fraction of full free volume required to provide polymer transition from the glassy into the rubbery state and the transition from crystalline phase to melt [12], comparatively enlarged  $T_g/T_m$  values in PVP–PEG blends imply facilitated PVP– PEG complex devitrification and PEG fusion. This finding is thoroughly explicable using the concept of fluctuation free volume formation under PVP–PEG mixing (Fig. 4). Being enriched with the fluctuation free volume, the PVP–PEG blends need less additional free volume for polymer transition into the viscoelastic state and for fusion.

While the  $T_g/T_m$  quantity varies linearly with blend composition, all the  $T_c$ -dependent quantities display the deviations from linearity at comparatively low PEG concentrations (Fig. 8). This fact is due to the contribution of molecular mobility to the  $T_c$  value as shown in the second paper of this series and illustrated below by the data in Fig. 13. The levelling off the response at high PEG weight fractions ( $w_{\text{PEG}} > 0.8$ ) for  $T_c / T_m$ and  $T_c/T_g$  plots (Fig. 8) results most likely from the disentanglement of PVP chains and gel–solution transition. Using optical microinterference measurements of the kinetics of high molecular weight PVP dissolution in liquid PEG-400, the PVP transition from swollen gel into solution has been recently found to occur at  $w<sub>PEG</sub> > 0.8$ .

Dependence of  $T_{g}/T_{m}$  quantity on PEG weight fraction in freshly prepared and stored blends with PVP is shown to obey the following equations:

$$
\frac{T_{\rm g}}{T_{\rm m}} = 0.82 - 0.06 w_{\rm PEG};
$$

 $R = -0.98$ ,  $p = 0.0006$  (fresh blends)

$$
\frac{T_{\rm g}}{T_{\rm m}} = 0.83 - 0.08 w_{\rm PEG}; R = -0.97,
$$

 $p = 0.0013$  (after blend exposure to water vapour for

three months at  $RH = 50\%$  and  $25^{\circ}$ C)

$$
\frac{T_{\rm g}}{T_{\rm m}} = 0.90 - 0.17 w_{\rm PEG}; R = -0.99,
$$

$$
p = 0.017 \text{ (dry blends)}
$$



Fig. 10. Dependence of  $T_c/T_m$  quantity on the PEG content in PVP–PEG blends freshly prepared and exposed to water vapour for three months at  $RH = 50\%$  and 25°C.

The storage conditions' effect on the  $T_g/T_m$  and  $T_c/T_m$ quantities are presented in Figs. 9 and 10. Providing appreciable contribution to phase behaviour, those nevertheless are not decisive for PVP–PEG compatibility and for interrelationships between phase transition temperatures.

The overall relationship between all the three transition temperatures in PVP–PEG compatible blends is given in Fig. 11. In the view of Eq. (9), the plot is indicative of the compositional behaviour of the  $\alpha_{\rm g}/\alpha_{\rm c}$  ratio in PVP–PEG blends. With the rise in PEG concentration, the  $\alpha_{\rm o}/\alpha_{\rm c}$ ratio increases linearly until a critical PEG concentration of  $w_{\text{PEG}} = 0.8$  is attained. Since only PEG is capable of forming a crystalline phase in PVP–PEG mixtures, the thermal expansion coefficient of crystalline polymer,  $\alpha_c$ , is a characteristic of PEG and is supposed to be approximately invariant with PEG content. Consequently, the increase in the  $\alpha_{g}/\alpha_{c}$  ratio with PEG weight fraction embeds, most likely, the increasing thermal expansion coefficient of PVP–PEG complex in glassy state,  $\alpha_{\rm g}$ .



Fig. 11. The ratio of  $T_m - T_g$  and  $T_c - T_g$  differences plotted against PEG concentration in blends with PVP.



Fig. 12. The Okui plot of  $T_c/T_m$  and  $T_g/T_m$  quantities in PVP–PEG blends of various prehistory, obtained with Eqs. (7)–(10).

Analysis of the  $T_g/T_m$  and  $T_c/T_m$  quantities for PVP–PEG compatible blends in terms of an Okui iso-volume state model and crystallization theory [1], featuring the activation energy of polymer segment migration through nucleus–melt interface,  $\Delta E$ , nucleation parameter,  $K$ , and a constant,  $C$ , defined by Eqs. (10) and (11), is presented in Fig. 12. The greater the *C* and  $\Delta E/K$  values the greater the  $T_g / T_m$ and  $T_c/T_m$  quantities, however the large variations in C and  $\Delta E/K$  in the region of their upper values cause relatively negligible changes in the  $T_g/T_m$  and  $T_c/T_m$ , as seen in Fig. 12. Apart from the minimum magnitudes of  $C = 1$  and  $\Delta E = 0$ , determined by definition (10), all the data points relating to the PVP–PEG blends fall in the region  $C > 5.5$ and  $\Delta E/K > 29$ . With increasing PEG content, the  $\Delta E/K$ gradually increases, tending to values as high as  $\Delta E/K =$ 525 at minimum PEG concentration when crystallization occurs ( $w_{\text{PEG}} = 0.53$ , Fig. 13). For individual homopolymers



Fig. 13. The dependence of  $\Delta E/K$  quantity, determined with Eqs. (10) and (11), on the PEG weight fraction in compatible blend with PVP.



Fig. 14. DSC traces upon heating of freshly prepared and aged PVP blends containing 52.9 wt% of PEG-400.

*C* has been found to vary from 3 to 9, corresponding to the change in  $\Delta E/K$  value between 8 and 80. The average value of  $\Delta E/K$  is about 23 *(C*  $\approx$  4.9) in most polymers [1].

Abnormally high  $\Delta E/K$  values for PEG-underloaded blends,  $(w<sub>PEG</sub> < 0.53)$ , give a plausible explanation for the lack of any crystallization and fusion processes observed in DSC heating traces of relevant mixtures (Fig. 1). It seems reasonable to say that the  $\Delta E/K \approx 500-600$  defines an upper limit of the range of activation energy, providing the existence of freezing PEG in the blends with PVP. As liquid PEG is added to glassy PVP, the blend  $T_g$  falls abruptly (Figs. 2 and 3). No crystalline phase is observed at the stage of PVP–PEG complexation (plasticization) (Fig. 2), while the  $T_g$  exceeds some critical magnitude. Upon the reduction of  $T_g$  value below this critical magnitude, the PEG crystallization becomes allowable. As is evident from Fig. 13, Eqs. (7), (8) and (10), comparatively negligible variations in  $T_g/T_m$  value within this critical region are due to dramatic changes in  $\Delta E/K$ , which enable or prohibit the crystalline phase formation.

This conjecture is conveniently illustrated in Fig. 14. In freshly prepared blends containing 52.9 wt% of PEG-400 and characterized with  $T_g = 214.4$  K; both PEG crystallization and fusion occurs upon sample heating through  $T_g$   $(T_c = 261.6 \text{ K}, T_m = 273 \text{ K},$  $T_p/T_m = 0.785$ ,  $\Delta E/K = 526$ ). As the same blend is allowed to relax for nine months at ambient temperatures and relative humidities, the  $T_g$  increases by as little as 5.3 K, and the freezing PEG ceases to be detected. The contribution of molecular mobility to the PEG crystallization in compatible blends with PVP is discussed in more detail in the second article of this series.

## **4. Conclusions**

DSC heating traces of cool-quenched PVP–PEG

compatible blends reveal a single, composition-dependent glass transition, intermediate between those of unblended polymer components, along with the phase transitions of excess PEG cold crystallization and melting, followed by the uppermost endotherm of sorbed water vaporization. In the compatible blends all the temperatures of phase transitions, as well as the  $T_g / T_m$ ,  $T_c / T_m$  and  $T_c / T_g$  quantities, are shown to be functions of the composition. Extremely large negative deviations from simple rules of mixing, expressed by the Fox and Gordon–Taylor equations, in the  $T<sub>g</sub>$ dependence on the composition are due to the earlier reported H-bonding of short-chain PEG terminal hydroxyl groups to carbonyls in the repeating units of PVP macromolecules, accompanied by free volume formation. The antagonistic contributions to the glass transition of conflicting strong favourable interaction between macromolecules of blended polymers and the free volume within PVP–PEG complex are embedded by a specific profile of the compositional dependence of a heat capacity change between the glassy and viscoelastic states,  $\Delta C_p$ . Occurrence of a crystalline phase in the compatible PVP–PEG blends is controlled by the molecular mobility of the crystallizable polymer.

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